Fourier Analysis 04-02-do24.
Review.
Consider the heat equation on the real line

$$\begin{cases} \frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}, & x \in \mathbb{R}, t > 0 \\ 1 & (x, 0) = f(x). \end{cases}$$
Using the method of Fourier transform, we have derived
a formal solution $U(x,t)$ as follows:
 $U(x,t) = f(x) + H_t(x),$
where
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where
 $\mathcal{U}(x,t) = \sqrt{4\pi t} e^{-\frac{x^2}{4t}}, & x \in \mathbb{R}, t > 0.$
We call $\{H_t\}_{t > 0}$ the Reat Rennel on \mathbb{R} .
Below we will provide a theoretic check.

Let
$$S(R)$$
 denote the collection of
all c^{∞} functions for R such that
for any $R, \ell \ge 0$,

$$\sup_{x \in \mathbb{R}} |x|^k \cdot |f^{(\ell)}(x)| < \infty$$

Here we collect some easily checked properties of S(R).

$$\forall f \in S(\mathbb{R}), f' \in S(\mathbb{R}).$$

•
$$\forall f \in S(\mathbb{R}), \quad x \neq \in S(\mathbb{R}).$$

More general, P(x) f(x) & S(iR) for each polynomial P.

Prop 2:
$$f \in S(R) \Leftrightarrow \hat{f} \in S(R)$$
.
Pf. We only prove the direction:
 $f \in S(R) \Rightarrow \hat{f} \in S(R)$.
The other direction follows from the Fourier inversion
formula.
Suppose $f \in S(R)$.
We need to show that
 $Sup |\hat{s}|^{k} \cdot |\hat{f}_{(3)}^{(\ell)}| < \omega \forall R, \ell \ge 0$.
 $\hat{s} \in R$
Notice that
 $(-2\pi i x)^{\ell} f(x) \xrightarrow{f} \hat{f}_{(3)}^{(\ell)}$

 $F(x) := \frac{d^{k} \left(\left(-2\pi i x \right)^{l} f(x) \right)}{dx^{k}} \xrightarrow{\mathcal{F}} \left(2\pi i \frac{k}{3} \right)^{k} \widehat{f}(\frac{k}{3})$ In particular $\sup_{\underline{x} \in \mathbb{R}} \left| (2\pi i \underline{x})^{k} \widehat{f}(\underline{x}) \right| \leq \int_{-\infty}^{\infty} \left| F(x) \right| dx$ < 10 (since FE S(IR)) $\Rightarrow \quad \sup_{3 \in \mathbb{R}} |\frac{3}{8}| \cdot |\frac{1}{5}| \cdot |\frac{1}{5}| < \infty$ 11

Thm 3. Let
$$f \in S(R)$$
 Let
 $U(x,t) = f * H_t(x)$.
Then
 $D \ U \in C^{\infty}(IR \times IR_t), \quad \frac{\partial U}{\partial t} = \frac{\partial U}{\partial x^2} \text{ on } IR \times IR_t$
 $E \ U(x,t) \Longrightarrow f(x) \quad as \quad t \to 0$
 $(3) \ \int_{IR} | \ U(x,t) - f(x) |^2 dx \to 0 \quad as \quad t \to 0$
 $Pf. \quad Since \ U = f * H_t, \quad both \ f, \ H_t \in S(R),$
it is not hard to show that
 $f * H_t \in S(IR) \quad \forall \quad t > 0.$
Now by Fourier inversion formula,
 $U(x,t) = \int_{-\infty}^{\infty} \widehat{U}(\underline{s},t) \ e^{2\pi i \cdot \underline{s} \times x}$

$$(*) = \int_{-\infty}^{\infty} \hat{f}(\hat{s}) \cdot e^{-4\pi^{2}\hat{s}^{2}t} e^{2\pi i \hat{s}x} d\hat{s}$$
Let us show
$$\frac{\partial u}{\partial t} = exists$$

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$$\frac{\partial u}{\partial t} = \int_{-\infty}^{\infty} \hat{f}(\hat{s}) \cdot \frac{e^{4\pi^{2}\hat{s}^{2}t} + \delta t}{\delta t} + \frac{4\pi^{2}\hat{s}^{2}t}{\delta t} + \frac{2\pi i \hat{s}x}{\delta t} + \frac{4\pi^{2}\hat{s}^{2}t}{\delta t} + \frac{2\pi i \hat{s}x}{\delta t} + \frac$$

By Lebesque's dominated Convergence Thm, $\lim_{\Delta t \to 0} \int_{-\infty}^{\infty} \frac{-4\pi^2 t}{f(t)} \frac{-4\pi^2 t}{2\pi} \frac{e^{-1}}{2\pi} \cdot \frac{2\pi}{3} x$ $= \int_{-\infty}^{\infty} f(\underline{s}) e^{-4\pi^{2}\underline{s}^{2}t} \cdot (-4\pi^{2}\underline{s}^{2}) \cdot e^{-d\underline{s}}$ By similarly arguments, we see that U∈ C[∞](R×R+). By (*), we have $\frac{\partial \dot{u}}{\partial \dot{x}} = \int_{-\infty}^{\infty} \hat{f}(\dot{z}) e^{-4\pi \dot{z}} \frac{\partial \dot{z}}{\partial \dot{z}} 2\pi \dot{z} \dot{z}$ and $\frac{\partial u}{\partial t} = \int_{-\infty}^{\infty} \hat{f}(\dot{z}) e^{-4\pi \dot{z}} \frac{\partial \dot{z}}{\partial \dot{z}} e^{2\pi i \dot{z}} \frac{\partial \dot{z}}{\partial \dot{z}}$ So $\frac{\partial U}{\partial t} = \frac{\partial \hat{U}}{\partial X^2}$ on $\mathbb{R} \times \mathbb{R}_+$. This proves (). Part (2) is a consequence that $\{H_t\}_{t>0}$ is a good kernel on \mathbb{R} .

Below we prove (3).
By Plancherel formula,

$$\int_{-\infty}^{\infty} | U(x,t) - f(x) |^{2} dx$$

$$= \int_{-\infty}^{\infty} | \hat{U}(\underline{3},t) - \hat{f}(\underline{3}) |^{2} d\underline{3}$$

$$= \int_{-\infty}^{\infty} | \hat{f}(\underline{3}) \cdot \underline{e}^{-4\pi^{2}\underline{3}^{2}t} - \hat{f}(\underline{3}) |^{2} d\underline{3}$$

$$= \int_{-\infty}^{\infty} | \hat{f}(\underline{3}) |^{2} \cdot | \underline{e}^{-4\pi^{2}\underline{3}^{2}t} - 1 |^{2} d\underline{3}$$
Notice that
$$\left| \hat{f}(\underline{3}) \right|^{2} \cdot | \underline{e}^{-4\pi^{2}\underline{3}^{2}t} - 1 |^{2} \leq \left| \hat{f}(\underline{3}) \right|^{2}$$

By DCT, $= \int_{\mathbb{R}} \lim_{t \to 0} |f(3)|^{2} |e^{-4\pi^{2}3^{2}t} - |d3|^{2}$ = 0.

Q: Is
$$U = \int_{x}^{x} \partial f_{t}$$
 the unique solution
to the heat equation

$$\int_{at}^{at} = \frac{\partial^{2} u}{\partial x^{2}} \quad \text{on } (\mathbb{R} \times \mathbb{R} + , \\ U(x, o) = \int_{x}^{\infty} (x, o) =$$